

# Technology of Lapped Optical-Fiber Couplers

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## Summary

The technological steps for the fabrication of lapped fiber couplers are described in detail. Choice of materials and of working parameters as well as the diagnostic measurements of the process are discussed. The guidelines of design are outlined and data on the performances of manufactured wavelength-division-multiplexing couplers are reported.

## 1 Introduction

Optical fiber couplers have been developed about ten years ago [1] and, since then, they have become well-established passive components of widespread use in communication networks and in sensors for the splitting and recombining of optical signals. In recent years, a large variety of special fiber couplers has been demonstrated, in which the splitting-ratio is either independent or selectively dependent on signal parameters such as wavelength, polarization, modal distribution, and (nonlinear) intensity.

As it is well known, there are three basic techniques currently used to fabricate fiber couplers, i.e.: *fused biconical*, *lapped* (or polished) and *micro-optic*. Several reviews [2-4] describe the fundamentals and compare the potentialities of each technique. However, let us briefly quote the commonly accepted advantages of the lapped-fiber technology as follows:

- (i) the alteration of the fiber structure is minimal, which allows to work easily on special fibers, such as high-birefringence fibers, for polarization splitting and polarization maintaining couplers, and also on dissimilar fibers;
- (ii) the open access to the interaction region allows the insertion of special intermediate layers (e.g.,  $\lambda$ -selective, birefringent or nonlinear) thus increasing the design flexibility;
- (iii) the inherent all-fiber structure has a negligible insertion loss.

On the other side, it shall be noted that environmentally related performances, ruggedness, volume production costs and yields of the lapped fiber technology are somewhere intermediate to those of the micro-optic and

fused biconical technologies. Moreover, the fabrication of multiport (in excess of  $2 \times 2$ ) couplers is difficult.

In this paper, we will focus on the fabrication of lapped fiber couplers, describing the choices of the abrasants and working plate, the optimization of the process parameters, the measurements for process characterization, and the assembly procedure. For the design of the coupler geometry (minimum separation and curvature radius), we will follow the well-known results of Snyder's [5] weak-coupling theory. As an example, we finally report some typical data on wavelength-division-multiplexing couplers fabricated by the described process.

## 2 Coupler geometry and fiber preparation

In a lapped coupler, the fibers are arranged as shown in Fig. 1. The fibers are mounted on a holder with a curvature radius  $R$  and the clad is removed by lapping down to a minimum separation  $d_0$  between the core axes.

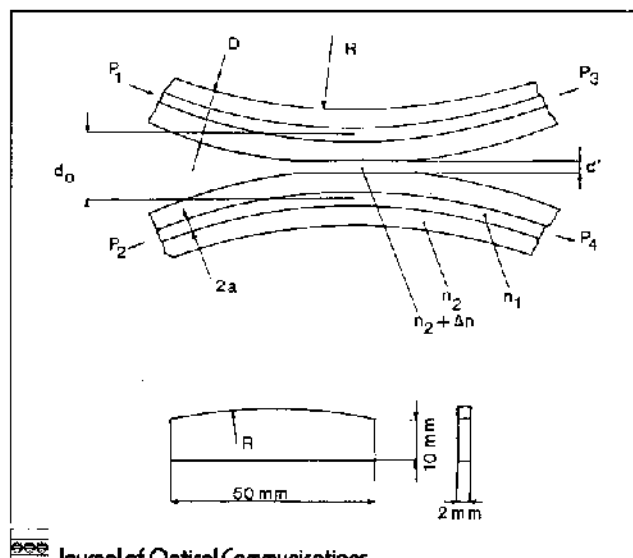


Fig. 1: Geometry of the coupler and of the fiber holder

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The most commonly reported geometry of the holder is that of the groove-in-a-block [5-7], particularly suited for fabricating couplers with an adjustable splitting-ratio. For fixed couplers, intended for assembly in a package, we preferred to emboss the fiber onto the edge of a curved plate (Fig. 1), so that the working time is considerably shortened since there is much less material to remove.

Another fast procedure for working the fiber without a holder is to lay it under tension on a polishing-wheel [8]; however, special care is necessary to avoid breakage of the thinned fibers during the assembling in the finished coupler.

The fiber is prepared for the lapping operation by stripping the coating (if any) with the help of a solvent (e.g. acetone) for a length equal to that of the holder, by aligning it on the centerline of the holder curved edge and cementing it with epoxy or other suitable resin, and finally by protecting it near the sharp holder edges with a drop of silicone resin.

The half-couplers are then ready for the lapping process described below. After this step, two half-couplers are mounted on a precision four-axis micropositioning stage and drawn together for the aligning procedure, using an index matching oil or glue between the fibers. The power outputs from ports 3 and 4 (Fig. 1) are monitored while launching first into port 1 and then into port 2. Small errors ( $\approx 1\mu\text{m}$ ) in the machining depth are compensated for in this phase by acting on the fiber relative displacement. When the desired splitting ratio is obtained, fast UV-curable resin is applied to cement the coupler.

Finally, the coupler is encapsulated in the package by means of silicone resin, with special care to avoid undesired asymmetrical shrinking during polymerization and the associated variation of the splitting ratio.

### 3 Lapping procedure and abrasant characterisation

To clamp in place the half-couplers and have a precisely lapped plane, in spite of their curved and blade-like shape, a fixture has been designed, consisting of two freely rotating coaxial cylindrical bodies, the inner one holding the half-coupler inside (Fig. 2). The external cylinder is secured to the moving arm of a standard rotating-plate lapping machine [9]. For all other details, the usual lapping practice [10] is employed.

Both alumina powders and diamond paste and slurry have been tested as abrasants. Routine inspections of surface finish were carried out with a metallographic optical microscope, sometimes backed by electron microscopy and light scattering analysis. Process parameters were monitored as well as other factors such as the frequency of fiber breakage and the uniformity of removal speed.

Alumina powder is dispersed in water and used on a rigid synthetic cloth disc. Among different compositions commercially available, alumina of high purity  $\alpha$ - $\gamma$  phase has given the best results in terms of reproducibility, and surface finish. After several trials with different particle sizes, it has been determined that the  $1\mu\text{m}$ -size alumina is fully adequate for a single step lapping, yielding a short

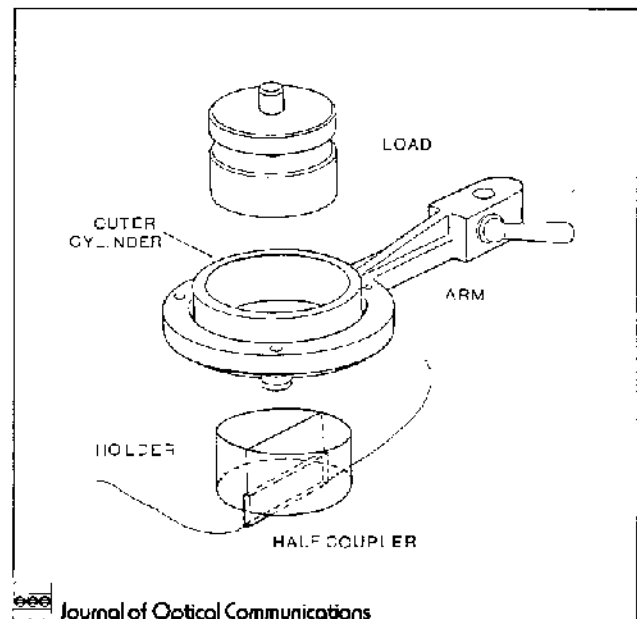


Fig. 2: Jig for lapping the half-couplers

( $\approx 10$  min.) machining time and yet a good finish quality even with no polishing.

It is interesting to correlate the removed depths  $s$  to the process variables, namely:

- (i) disc speed  $v$  (turns/min);
- (ii) pressure of the arm  $p$  ( $\text{kg}/\text{cm}^2$ );
- (iii) alumina particle size  $\phi$  ( $\mu\text{m}$ ), and
- (iv) working time  $t$  (hours).

Using flat samples of constant area ( $20\text{ mm}^2$ ), we obtained the curves of  $s$  versus  $t$  reported in Fig. 3. The following empirical relation is found to fit well the experimental data:

$$s [\mu\text{m}] = h p^\alpha \phi^\beta v^\gamma t [\text{hours}] \quad (1)$$

where the constants are given by

$$h = 1.2 \quad \alpha = 1.0 \quad \beta = 0.45 \quad \gamma = 1.0. \quad (2)$$

These values are the result of a best fit around:  $v = 100$  turns/min,  $p = 1\text{ kg}/\text{cm}^2$ ,  $\phi = 1\mu\text{m}$ , and apply in the range  $v = 30\text{--}250$  turns/min,  $p = 0.5\text{--}5\text{ kg}/\text{cm}^2$ ,  $\phi = 0.3\text{--}5\mu\text{m}$ .

For comparison, we also report the values of the same parameters for a diamond paste ( $1\mu\text{m}$ -size) on lead-charged resin disc, around the same values of  $v$ ,  $p$ , i.e.:

$$h = 3.8 \quad \alpha = 1.0 \quad \beta = 1.25 \quad \gamma = 0.5. \quad (3)$$

The differences between (2) and (3) can be attributed not only to the different hardness of the abrasants, but also to the different machining mechanism. Diamond particles are likely to be embedded into the disc, thus, exerting a true lapping action, while alumina moves loosely on the cloth and its action is closer to polishing.

With the  $1\mu\text{m}$ -size diamond paste, a larger number of fiber breakages and some evidence of under-surface

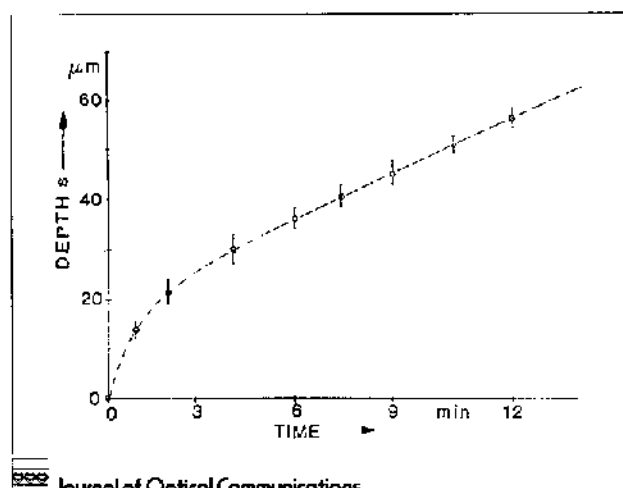


Fig. 3: Removed depth vs time for flat glass samples (20 mm<sup>2</sup> area) lapped with alumina (1 μm particle size) at different values of arm pressure  $p$  and plate speed  $v$

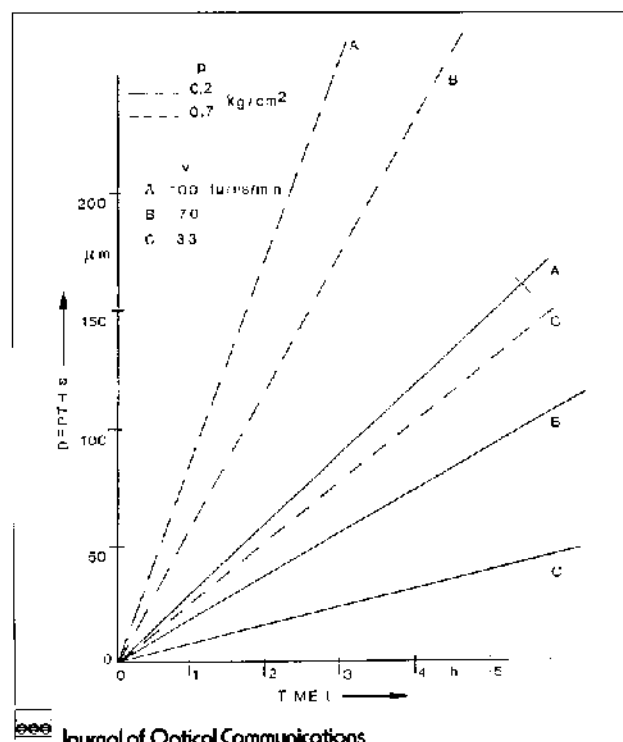


Fig. 4: Removed depth vs time for a silica fiber lapped at constant weight ( $W = 300$  g) and speed ( $v = 33$  turns/min) with 1 μm alumina powder

damage [1] are found, while finer sizes yield too long machining times.

For the lapping of fibers mounted on half-couplers with 1 μm-size alumina, we plot in Fig. 4 the removed depth  $s$  versus the working time  $t$ . The diagram is not linear at short times because the contact area between plate and fiber increases rapidly as lapping progresses, while the load on the tool and the disc speed are constant. At longer times, the regime becomes quasi-linear, and (1) holds with the following values of parameters and process variables:

$$h = 1.2, \quad \alpha = 0.3, \quad \beta = 0.45, \quad \gamma = 1.0, \\ v = 30\text{--}250 \text{ turns/min}, \quad p = 10\text{--}50 \text{ kg/cm}^2, \\ \phi = 0.3\text{--}5 \mu\text{m}.$$

From the spread of the experimental data (Fig. 3), the working time allows to estimate the removed depth with an accuracy of 2–5 μm.

#### 4 Process control and diagnostic measurements

An accurate control of the distance from the core/clad boundary is necessary in working the half-couplers, because the coupling factor varies appreciably as the minimum distance  $d_0$  is changed on a 1 μm-scale for low-order coupling [2–5], and even more rapidly for high-order coupling. In addition, one shall take into account the tolerance on the fiber outer diameter, which is often well in excess of 1 μm in ordinary telecommunication fibers, and is usually found uncorrelated from sample to sample a few meters apart along the fiber.

A short cut to avoid this problem is obviously to overwork the fiber down to approximately the core/clad boundary, and insert between the two half-couplers a thick layer of index-matching resin ( $d'$  in Fig. 1), squeezed to the appropriate thickness during the trimming operation. However, since the index of refraction varies during the resin polymerization and with aging, the trimming operation becomes more complicate and the long-time stability is objectionable; also, an unavoidable mismatch is expected versus wavelength. Therefore, only a relatively thin (of the order of 1 μm) index-matching layer is advisable.

Two kinds of diagnostics are required: *process* diagnostics, to be assessed once and necessary to establish the dependence of the removed depth on the process parameters, and *routine* diagnostics, necessary during the working of the half-couplers for lapping the fiber down to the design depth.

A destructive method has been used for the process diagnostics: the fiber is detached from its holder by a solvent and its thickness is measured under a metallographic microscope using a precision graticule scale, obtaining readings with a typical accuracy of  $\approx 0.5$  μm.

Other nondestructive diagnostics based on the measurement of the removed depth are unfortunately much less accurate and, in addition, they are affected by the outer diameter error. As said above, the calculation of the removed depth from the elapsed time yields an unacceptable error (in the range 2–5 μm), and this approach can then be used only as a first-step estimate, to rapidly remove most of the clad down to a safe distance from the fiber core. Another method, i.e., the correlation between the removed depth and the dimensions (axis-length) of the flat, ellipse-like section created on the fiber by lapping, is even less accurate (typ. 10 μm) because of the difficult visual identification of the ill-defined ellipse borderline.

More effective are those routines aimed to measure directly the actual distance of the lapped surface from the core. Among them, the technique of launching power in the half-coupler and measuring the transmitted power attenuation during lapping, originally proposed for diamond polishing [10], has been found difficult to implement with loose alumina on a cloth disc, because the fiber contact to the disc varies as the machine arm moves back and forth, and thus the output power exhibits strong fluctuations. To get reliable results, this method requires a

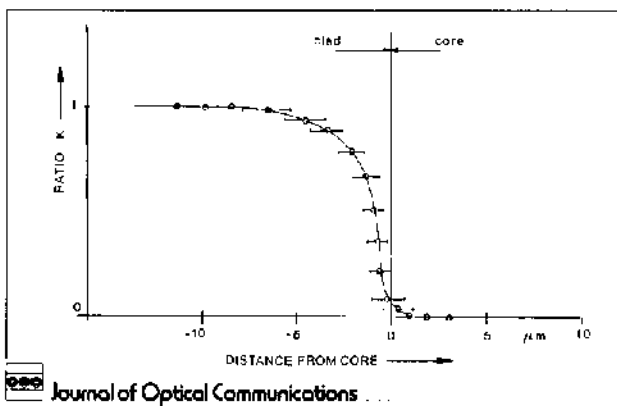


Fig. 5: Attenuation  $K$  vs distance due to an absorbing tape applied onto the lapped fiber; the error bars include the concentricity error of the core

stable and suitable refractive index of the medium surrounding the fiber being lapped, which interferes with the preparation of the abradant suspension.

Instead, a precise but time-consuming method is that of measuring the power extracted from the lapped fiber by oils of different refractive indexes [13]. This method has been estimated accurate down to  $0.5\text{--}1\ \mu\text{m}$ . We use a modified version as a procedure sufficiently fast for routine working. With a rough estimate based on the elapsed working-time, the core is approached to within  $5\ \mu\text{m}$ , lapping is stopped and the half coupler is washed and dried. Power is launched at one fiber end while the transmitted power  $P$  is measured at the other end. Then, an absorbing tape (a common tape with anaerobic glue) is pressed onto the lapped area to absorb a well-defined and reproducible fraction of the power available for coupling, and the new power value  $P'$  is measured. The ratio  $K = P'/P$  is very sensitive to the distance from the core in the region around the core-clad interface, as shown in Fig. 5, and it gives a typical accuracy better than  $1\ \mu\text{m}$  on a useful working range in excess of  $5\ \mu\text{m}$  from the core boundary. The result of Fig. 5 is obtained by using a  $633\ \text{nm}$  He-Ne laser source and a  $1300\ \text{nm}$  monomode fiber (see Sect. 6). Because of the dependence on the normalized frequency  $v$  of the fiber, the working range can be extended using longer wavelengths (e.g. to about  $8\ \mu\text{m}$  with a  $850\ \text{nm}$  diode laser), with a corresponding decrease in accuracy.

## 5 Design rules

In the lapped-coupler geometry (Fig. 1), the design parameters to be selected are the minimum separation  $d_0$  and the radius of curvature  $R$ . Restricting ourselves to couplers made with identical monomode fibers, and neglecting for the moment the effect of the intermediate layer ( $d' = 0$ ), we apply the basic results of Snyder's low perturbation theory, which represents a good approximation in the weakly guiding regime [5, 14]. Our aim is to evaluate the typical values of the design parameters for different kinds of couplers, namely: fixed-ratio, WDM, mode- and polarization-splitters.

For an ideal lossless coupler, if  $P_1$  is the power launched into port 1, the powers  $P_3$  and  $P_4$  at the direct and crossed ports (Fig. 1) are given by [4]:

$$P_3/P_1 = \eta = \sin^2(\pi L/2L_c), \quad (4)$$

$$P_4/P_1 = 1 - \eta = \cos^2(\pi L/2L_c). \quad (5)$$

In these equations,  $L_c$  is the coupling length between the symmetrical and antisymmetrical modes of the two-core structure, and  $L$  is the effective interaction length; these quantities can be expressed in terms of  $d_0$ ,  $R$ , and of the fiber geometrical parameters and guiding constants  $v$ ,  $u$ ,  $w$  as [5]:

$$L_c = \pi/2 C_0 = \pi/[2(\pi \delta/w d a)^{1/2} (u^2/v^3) \exp(-w d_0/a)/K_1^2(w)] \quad (6)$$

$$L = (\pi R a/w)^{1/2} \quad (7)$$

where  $C_0$  is the coupling coefficient evaluated at the minimum core distance  $d_0$ , and  $\delta = 1 - n_2/n_1$ .

Power-division couplers exploit the dependence of  $L_c$  on  $R$  and  $d_0$ . The coupler is usually tuned on the first order of coupling for the best stability. A weak dependence of the splitting ratio on  $\lambda$  is often desirable and can be obtained either by using dissimilar fibers or by properly selecting the curvature radius  $R$ , while the splitting ratio is trimmed by acting on  $d_0$  [15].

Typical values of the effective interaction length  $L$  are of the order of  $1\ \text{mm}$ , while the length of the lapped region  $L$  (the major axis of the ellipse worked on the fiber) is about  $10\ \text{mm}$ , a value well in the reach of the lapping technology.

Couplers for the multiplexing of wavelength, polarization state or mode-order exploit the dependence of  $L_c$  on the guiding constants  $u$ ,  $v$  and  $w$ . On the diagram of  $\eta$  vs.  $L$ , solutions are found when two curves have a maximum and a minimum of transmission in correspondence. This requires an effective interaction length  $L^*$  given by:

$$L^* \cong L_{c1} L_{c2} / (L_{c1} - L_{c2}). \quad (8)$$

If the coupling lengths  $L_{c1}$  and  $L_{c2}$  are very close, the interaction length can be much larger than  $L_{c1}$  and  $L_{c2}$ , i.e., the first order of multiplexing corresponds to a high order of power-splitting coupling. Using in (9) the approximate expression  $w \cong gv - 1$ , where  $g \cong 1.14$  [16],  $L^*$  is found as:

$$L^* \cong (L_{c1} + L_{c2}) / \{g \Delta v [(1/(gv - 1) - 6/gv) + (gv - 4)/(1 + 0.25 g^2 v^2 - 2gv)]\}. \quad (9)$$

Typical values of  $L^*$  are in the range of a few mm for a wavelength multiplexing of the second and third windows, using a standard monomode fiber; also, the same order of magnitude is found for polarization multiplexers made with high-birefringence fibers.

Expressions for mode multiplexing are considerably more complicated, even in the simple case of  $LP_{01}$  and  $LP_{11}$  splitting. Interaction lengths of a few mm have been reported [17] for mode multiplexers based on dissimilar fibers; we have found this result consistent with a numerical evaluations and several experimental trials.

Nonlinear couplers exploit the dependence of the refractive index of a thin layer of nonlinear material, inserted between two standard half-couplers, on the intensity of the propagating wave. For a coupler structure with intermediate layer, it is found [12] that (4-7) still hold, provided that the coupling coefficient  $C_0$  is multiplied by

a correction factor  $1+h$ . If  $d'$  and  $n_{nl} = n_2 + \Delta n$  are the thickness and the refractive index of the nonlinear layer, respectively (see Fig. 1), in the assumption of weak perturbation one can find [12]:

$$h = vd'(n_{nl} - n_2) \pi / (w\lambda \sqrt{2\delta}). \quad (10)$$

The nonlinear refractive index can be expressed in terms of the irradiance  $\Psi$  ( $\text{W}/\text{cm}^2$ ) as:  $n_{nl} = n_0 + \gamma\Psi$ , where the third-order susceptibility  $\gamma$  is very low in silica (e.g.,  $\approx 10^{-10} \text{ cm}^2/\text{W}$ ) [18], but can reach  $10^{-3} - 10^{-4} \text{ cm}^2/\text{W}$  in doped glasses and GaAs. A coupler designed for maximum transmission between ports 1 and 4 at low power will then switch to port 3 when power is increased beyond a specific level. For the coupler described in Section 6, with  $L = 3 \text{ mm}$  and  $d_0 = 9.1 \mu\text{m}$ , switching is obtained for a variation of  $d' \Delta n$  of the order of  $0.01 \mu\text{m}$ . Using a suitable nonlinear glass, such a variation can be obtained with an input power ranging from a few hundred to a few thousand watts.

### 6 WDM couplers

As an illustrative example of fabricated lapped-fiber couplers, we report some results on wavelength-division-multiplexing (WDM) couplers intended for the splitting at 1300 and 1550 nm. A standard monomode telecommunication fiber (Pirelli SM01) has been used, whose relevant parameters are:  $\lambda_{\text{cutoff}} = 1230 \text{ nm}$ , core radius  $a = 4.55 \mu\text{m}$ , clad diameter  $D = 125 \mu\text{m}$ ;  $n_2 = 1.4514$ ,  $n_1 = 1.4458$ , (at 1300 nm). From numerical calculations of (4-7), the coupling factor  $\eta$  is evaluated as a function of the wavelength for several values of the parameters  $d_0$  and  $R$ . Below a certain value of the curvature radius ( $R < 0.8 \text{ m}$ ), complete wavelength separation is not achieved, while several solutions for  $d_0$  and  $R$  are found for  $R > 0.9 \text{ m}$ . However, at higher orders of coupling, the control of lapping and trimming becomes critical as  $\eta$  varies rapidly with  $d_0$ . A typical plot of  $\eta$  versus  $d_0$  is reported in Fig. 6 for  $R = 1 \text{ m}$  (full lines), showing two possible choices of  $d_0$  for the multiplexing between  $\lambda = 1300$  and  $1550 \text{ nm}$ . Also plotted in Fig. 6 (dotted lines) is the effective  $\eta$  when the correction due to the matching layer between the fibres is taken into account ( $\Delta n = -0.01$ ). Both 1300 and 1550 nm curves are shifted, but their trend is essentially unchanged and it can be compensated for in the final trimming of the coupler. The calculated wavelength dependence of the coupling factor  $\eta$  is plotted in Fig. 7, for  $R = 1 \text{ m}$  and  $d_0 = 9.9 \mu\text{m}$  (right-hand side arrows of Fig. 6).

Experimental points relative to a typical WDM coupler, representative of the performances obtained in a batch fabricated by lapping, are also reported in Fig. 7. The coupler was trimmed to match the working wavelengths, compensating for the effect of the intermediate layer and for machining errors. The spectral transmission response of the same device is reported in Fig. 8. The excess loss of the WDM coupler is 1 dB and the extinction ratio is 23 dB. Other average performances of the batch were: back isolation  $P_2/P_1 < 35 \text{ dB}$ , error on nominal multiplexing wavelengths  $\lambda_1, \lambda_2: < 10 \text{ nm}$ , error in wavelength separation  $\Delta\lambda = \lambda_2 - \lambda_1 < 20 \text{ nm}$ .

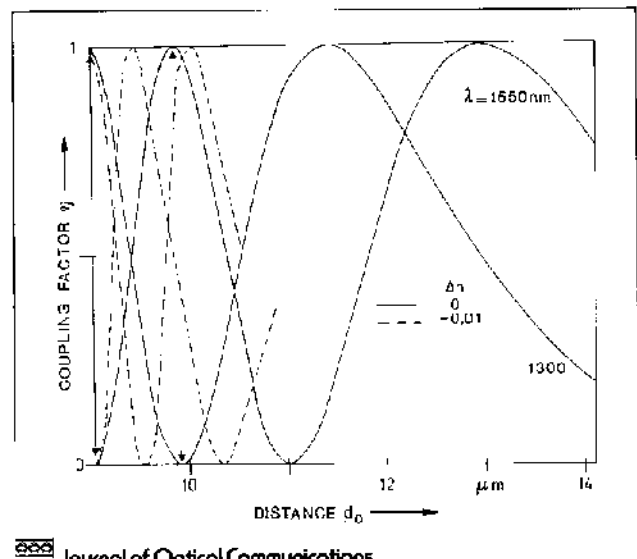


Fig. 6: Calculated coupling factor  $\eta$  versus core distance  $d_0$ , for  $R = 1 \text{ m}$  (full lines); multiplexing between 1300 and 1550 nm is indicated by the arrows; dotted lines show the correction for a  $\Delta n = -0.01$  matching layer ( $d' = 0$  at  $d_0 = 9.1 \mu\text{m}$ )

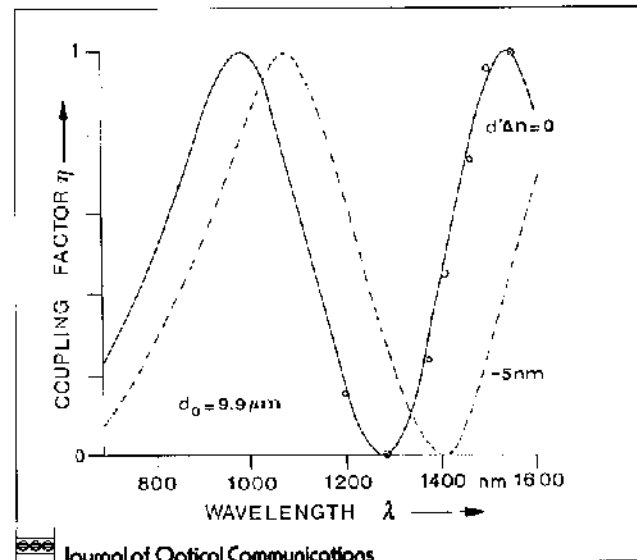


Fig. 7: Calculated coupling factor  $\eta$  versus wavelength  $\lambda$  (with  $R = 1 \text{ m}$ ,  $d_0 = 9.9 \mu\text{m}$ ) for the WDM coupler with no intermediate layer (full lines) and with a thin intermediate layer  $d' \Delta n = -5 \text{ nm}$  (dotted lines); experimental points are relative to the WDM coupler characterized in Fig. 8

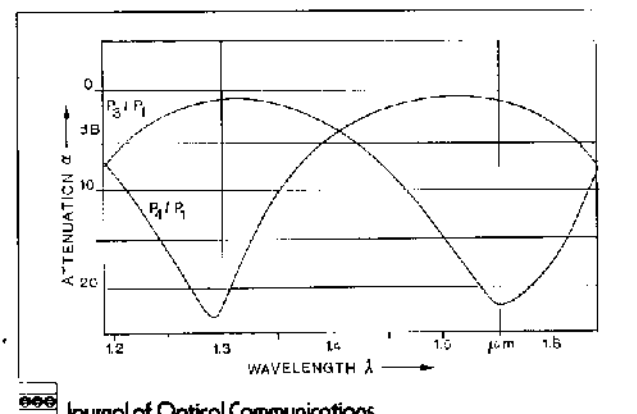


Fig. 8: Experimental result representative of a lapped WDM coupler: direct ( $P_3/P_1$ ) and crossed ( $P_4/P_1$ ) coupling factors versus wavelength

## 7 Acknowledgement

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