

## Thick-film piezo actuators for SOS devices

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**Abstract** - We report on a new method for pathlength actuation of silica-on-silicon (SOS) integrated optics devices. Compared to the commonly employed thermal actuation by dissipation in a resistive film, the proposed method has the advantages of no dc power consumption and of a faster (down to microseconds) response. Also, we discuss the integrated fabrication of several elements by means of thick-film silk-screen printing of piezoelectric inks and appropriate conductive inks for the electrodes.

### I INTRODUCTION

Virtually all integrated-optic devices require a sort of trimming of pathlength to work properly. Examples are interferometric devices like AWG's, Mach-Zehnder modulators, but also couplers, switches, etc. The pathlength actuation is commonly achieved by a thin resistive layer evaporated over the waveguide, in which by Joule-dissipation one gets a thermo-optical control of pathlength.

Thermal actuation is rather slow (response times 0.1-1 ms) and a substantial power is demanded in CW operation, typically 300-500 mW per element, thus impeding the use of extensive thermal control on chip. In this paper we propose a piezo actuation [1] and show that it is feasible for integrated fabrication on a chip, with an improvement of both the power dissipation (zero in CW) and the speed of response (down to microseconds).

In the following we compare thermal and piezo actuation finding a substantial agreement with data measured on routinely fabricated SOS Mach-Zehnders (made by 10-mm waveguides and two 50-50 2x2 couplers). For the comparison, the interferometers were actuated on one leg either thermally or through a 7-mm length PZT-piezo cemented to the substrate (fig.1).

### II ANALYSIS

An analysis of thermal and piezo actuations has been carried out in [1], and the main results are summarized below. The pathlength variation in the guide, due to thermal dissipation of a power  $P$ , can be written as:

$$\Delta(nL) = \zeta nL \Delta T_g = \zeta nL K_t P \eta \quad (1)$$

where  $L$  is the actuator length,  $K_t$  ( $^{\circ}\text{C}/\text{W}$ ) is the thermal resistance from the layer body to the ambient,  $\zeta = dn/dT + dL/LdT$  is the thermo-optical coefficient of the guide ( $\zeta = 8.1 \cdot 10^{-6} \text{ }^{\circ}\text{C}^{-1}$  in silica [2]) and  $\eta$  is an efficiency factor.

Expliciting the thermal resistance as [3]:  $K_t = [\ln(4L/W)]/\pi\kappa L$ , where  $\kappa = 0.014 \text{ Wcm}^{-1} \text{ }^{\circ}\text{C}^{-1}$  is the thermal conductivity of silica, we get:

$$\Delta(nL) = \zeta n [\ln(4L/W)/\pi\kappa] P \eta \quad (2)$$

and, by letting  $\Delta(nL) = \lambda/2$  as the switching condition, we get for the drive power:

$$P_{\lambda/2} = \pi\kappa\lambda/[2n \zeta \cdot \ln(4L/W)\eta] \quad (3)$$

For our experiment at  $\lambda = 1500 \text{ nm}$ , with  $L = 4.2 \text{ mm}$ ,  $W = 25 \text{ } \mu\text{m}$ ,  $\eta = 0.12$  (from simulations), eq.(3) gives  $P_{\lambda/2} = 370 \text{ mW}$ , a value close to the measured one (420 mW). Such a high value is primarily due to the small constant  $\zeta$  of silica [eq.(3)]. About response time, the main effect is the time constant  $\tau = K_t \cdot C_t$  of the layer,  $C_t$  being the thermal capacity of the layer. As it is [3]:  $C_t = c_{sp} \pi W^2 L \cdot \ln(4L/W)$ , where  $c_{sp} = 0.24 \text{ J cm}^{-3} \text{ }^{\circ}\text{C}^{-1}$  is the specific heat of silica, we have found [1]:

$$\tau = (c_{sp}/\kappa) [W \cdot \ln(4L/W)]^2 \quad (4)$$

Eq.(4) yields  $\tau = 0.45 \text{ ms}$ , whence a risetime (10-90% points)  $T_r = 0.99 \text{ ms}$  in fair agreement with the experimentally measured ( $T_r = 0.75 \text{ ms}$ ).

Now, going to analyze the piezoelectric actuation, the pathlength variation is [1]:

$$\Delta(nL) = p_{11} nL S_1 = p_{11} nL d_{13} E_3 \eta \quad (5)$$

where  $S_1$  is the strain,  $p_{11}$  is the strain-optical coefficient of the guide ( $= 0.121$  in silica), and  $d_{13}$  is the piezoelectric coefficient of the PZT ceramic ( $d_{13} = -420 \cdot 10^{-12} \text{ C/N}$ ); the factor  $\eta$  again accounts for incomplete stress coupling from piezo to waveguide. Letting  $\Delta(nL) = \lambda/2$  in (5) gives the switching field as:

$$E_{\lambda/2} = (\lambda/2 p_{11} d_{13} nL)/\eta \quad (6)$$

and for  $L = 7 \text{ mm}$  as in our device it is  $E_{\lambda/2} = 1500 \text{ V/mm}$ , or  $V_{\lambda/2} = 1500 \text{ V}$  for the switching voltage that corresponds to  $\eta = 0.44$  in the experiment.

The response time is determined [4] by the time spent, at the velocity of sound  $v_{ac}$  ( $= 3000 \text{ m/s}$ ), to go across the sample thickness  $T$  and establish a standing

wave. Thus we have  $Q$  (quality factor of the ceramic) times the period of the resonant frequency  $v_{ac}/2T$  as the minimum response time:

$$\tau = 2QT/v_{ac} \quad (7)$$

For  $T=1\text{mm}$ ,  $Q=2$ , (7) gives  $\tau=1.2 \mu\text{s}$ , near to measured data ( $2 \mu\text{s}$ ). The product  $\tau E\lambda/2$  is independent from  $L$  and one can trade-off half wave voltage and speed. In the piezo actuation, no power is dissipated in dc. On switchings, however, the energy  $1/2CV^2$  ( $1.4\text{mJ}$  in our setup) is dissipated in the external circuit.

### III INTEGRATED PIEZOS

The drive voltage can be substantially decreased in a practical design, e.g., by employing two piezos in push-pull arrangement to halve the voltage. Also, thinner elements or a multilayer ceramic could be used, reducing by a factor 3-5 the required voltage, down to a reasonable 100-200V.

An approach [5] easy to implement is shown in fig.2: two multilayer SMD ceramics, much alike SMD (surface-mount-device) high- $\epsilon$  capacitors, are used in push-pull to drive the guides. The SMD's can be mounted with pick-and-place on solder pad previously deposited on the SOS with standard electronic manufacturing assembly techniques, very common in the thick-film circuit technology.

With a multilayer piezo ( $n=8$  layers) of  $L=5\text{mm}$ ,  $W=1\text{mm}$ ,  $T=2\text{mm}$ , a drive voltage of 50 V and a switching time of  $2\mu\text{s}$  are anticipated. With this small  $W$ , several actuators can be placed aside for actuation in multiguide SOS.

Even better to increase the integration density is the approach of thick-film printing, going down to a few mil line resolution as required, e.g., by AWG's (fig.3). Piezoelectric inks have been reported [6] for sensor applications and they exhibit, when fired with a thermal profile similar to that of thick film electronic devices, piezoelectric constants close to the bulk values, and good long-term stability.

### REFERENCES

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[6] B. Morten et al.: PZT-based Thick Films for Piezoelectric Pressure Sensor, Hybrid Circuits, 28 may 1992, pp.25-28.

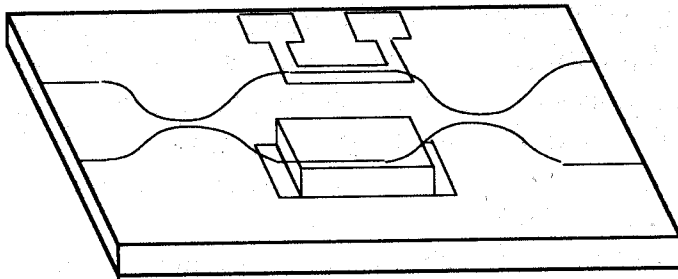


Fig.1. A SOS Mach-Zehnder interferometer with thermal and piezo actuators on legs

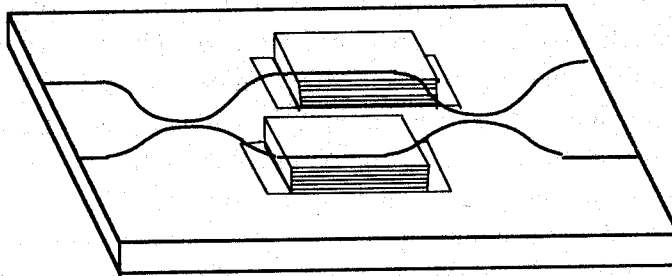


Fig.2. Actuation by means of SMD piezos

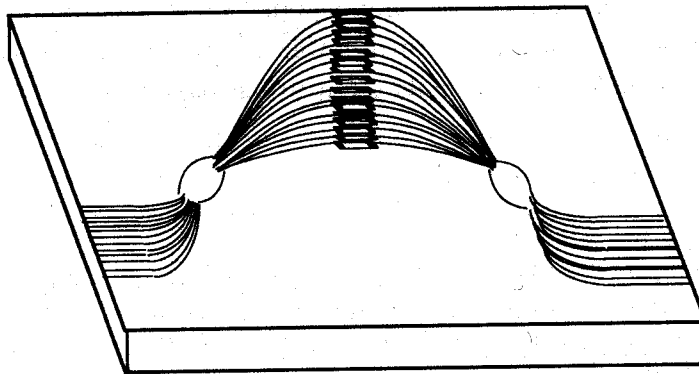


Fig.3. Actuation by means of thick-film ceramic-ink piezos